

Home Search Collections Journals About Contact us My IOPscience

A new Q-matrix in the eight-vertex model

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

2007 J. Phys. A: Math. Theor. 40 4075

(http://iopscience.iop.org/1751-8121/40/15/002)

View the table of contents for this issue, or go to the journal homepage for more

Download details: IP Address: 171.66.16.108 The article was downloaded on 03/06/2010 at 05:06

Please note that terms and conditions apply.

J. Phys. A: Math. Theor. 40 (2007) 4075-4086

doi:10.1088/1751-8113/40/15/002

4075

A new *Q*-matrix in the eight-vertex model

Klaus Fabricius

Physics Department, University of Wuppertal, 42097 Wuppertal, Germany

E-mail: Fabricius@theorie.physik.uni-wuppertal.de

Received 11 December 2006, in final form 27 February 2007 Published 23 March 2007 Online at stacks.iop.org/JPhysA/40/4075

Abstract

We construct a *Q*-matrix for the eight-vertex model at roots of unity for crossing parameter $\eta = 2mK/L$ with odd *L*, a case for which the existing constructions do not work. The new *Q*-matrix \hat{Q} depends on the spectral parameter *v* and also on a free parameter *t*. For t = 0, \hat{Q} has the standard properties. For $t \neq 0$, however, it does not commute with the operator *S* nor with itself for different values of the spectral parameter. We show that the six-vertex limit of $\hat{Q}(v, t = iK'/2)$ exists.

PACS number: 75.10.Jm

An essential tool in Baxter's solution of the eight-vertex model [1-4] is the *Q*-matrix which satisfies the TQ equation

$$T(v)Q(v) = [\rho h(v - \eta)]^{N}Q(v + 2\eta) + [\rho h(v + \eta)]^{N}Q(v - 2\eta)$$
(1)

and commutes with *T*. Here T(v) is the transfer matrix of the eight-vertex model (A.1). Combined with periodicity properties of Q(v) in the complex *v*-plane equation (1) leads to the derivation of Bethe's equations and the solution of the model. For generic values of the crossing parameter η the transfer matrix *T* has a non-degenerate spectrum. For rational values of η/K however this is not the case. This leads to the existence of different *Q*-matrices which all satisfy equation (1). In [1], Baxter constructs a *Q*-matrix valid for

$$2L\eta = 2m_1K + \mathrm{i}m_2K' \tag{2}$$

with integers m_1, m_2, L . In [2], Baxter derived a *Q*-matrix valid for generic values of η . As these *Q*-matrices are different we distinguish them by writing Q_{72} and Q_{73} respectively for the constructions in [1, 2]. It turned out, however, that Q_{72} has interesting properties beyond its role in equation (1) *because* of its restriction to rational values of η/K .

In [5], it is conjectured that $Q_{72}(v)$ satisfies the following functional relation.

For N even and $\eta = m_1 K/L$ where either L even or L and m_1 odd

$$e^{-N\pi i v/2K} Q_{72}(v - iK') = A \sum_{l=0}^{L-1} h^N (v - (2l+1)\eta) \frac{Q_{72}(v)}{Q_{72}(v - 2l\eta)Q_{72}(v - 2(l+1)\eta)}, \quad (3)$$

1751-8113/07/154075+12\$30.00 © 2007 IOP Publishing Ltd Printed in the UK

A is a normalizing constant matrix independent of v that commutes with Q_{72} and $h(v) = H(v)\Theta(v)$. There is a proof of this conjecture valid for L = 2 in [6]. This functional relation is important as it allows the conclusion that the dimension of eigenspaces of degenerate eigenvalues of the *T*-matrix is a power of 2, a result also true in the six-vertex model provided the roots of the Drinfeld polynomial of the loop algebra symmetry are distinct [7].

The reason why the case L odd and m_1 even is excluded in (3) is that Q_{72} does not exist in this case [5].

The purpose of this paper is to close this gap¹. We construct for even N a Q-matrix which exists for $\eta = 2mK/L$ for odd L which satisfies the functional relation (3). Beyond that we shall show that for $\eta = 2mK/L$ a more general Q-matrix exists which depends on a free parameter t and which does not commute with R and S where

$$R = \underbrace{\sigma_1 \otimes \sigma_1 \otimes \cdots \otimes \sigma_1}_{\text{N factors}} \qquad S = \underbrace{\sigma_3 \otimes \sigma_3 \otimes \cdots \otimes \sigma_3}_{\text{N factors}} \tag{4}$$

and not even with itself for different spectral parameters

$$[Q(v_1, t), Q(v_2, t)] \neq 0.$$

This phenomenon has also been observed by Bazhanov and Stroganov in [8] for their columnto-column transfer matrix T_{col} which acts like a *Q*-matrix in the six-vertex model: it satisfies (1) and it commutes with T_6 . But it does not commute with itself for different arguments.

We use the notation of Baxter's 1972 paper. We denote our new *Q*-operators by \hat{Q}_R , \hat{Q}_L and \hat{Q} . They depend on two arguments *v* and *t*, e.g. $\hat{Q}_R(v, t)$. For brevity we shall write $\hat{Q}_R(v)$ instead of $\hat{Q}_R(v, 0)$. The symbols *Q*, Q_R etc refer to all types of *Q* matrices.

The plan of this paper is as follows. In section 1 we describe the various steps in the construction of Q. We first outline in section 1.1 the general method developed by Baxter and his solution leading to Q_{72} . In section 1.2 we present our new \hat{Q}_R operator and describe its range of validity. In section 1.3 we introduce the matrix Q_L and show in section 1.4 that the famous equation $Q_L(u)Q_R(v) = Q_L(v)Q_R(u)$ which Baxter proved for Q_{72} and Q_{73} is also satisfied by $\hat{Q}(v)$. In section 2 we study the quasiperiodicity properties of Q(v) and show that there exists a link between quasiperiodicity of Q_R with quasiperiod iK' and non-existence of Q_R^{-1} . We summarize in section 3 the properties of $\hat{Q}(v)$ and describe in section 4 the exotic properties of $\hat{Q}(v, t)$ for $t \neq 0$.

1. Construction of a *Q*-matrix for $\eta = 2mK/L$

1.1. Baxter's construction of Q_{72}

The goal is to find a matrix Q_R of the form

$$[Q_R(v)]_{\alpha|\beta} = \operatorname{Tr} S_R(\alpha_1, \beta_1) S_R(\alpha_2, \beta_2) \cdots S_R(\alpha_N, \beta_N),$$
(5)

where α_i and $\beta_i = \pm 1$ and $S_R(\alpha, \beta)$ is a matrix of size $L \times L$ such that Q_R satisfies

$$T(v)Q_R(v) = [\rho h(v-\eta)]^N Q_R(v+2\eta) + [\rho h(v+\eta)]^N Q_R(v-2\eta).$$
(6)

The Q-matrix occurring in equation (1) is then

$$Q(v) = Q_R(v)Q_R^{-1}(v_0)$$
(7)

for some constant v_0 . Therefore, it is necessary that $Q_R(v)$ is a regular matrix. The problem to construct a Q_R of the form (5) satisfying (6) is posed and solved by Baxter in appendix C

¹ There now exists a related investigation by Roan [11].

of [1]. In order to construct a Q_R -matrix which is regular for $\eta = mK/L$ for even *m* and odd *L*, we shall search for other solutions of Baxter's fundamental equations.

These equations are (see (C10), (C11) in [1])

$$(ap_n - bp_m)S_R(+, \beta)_{m,n} + (d - cp_m p_n)S_R(-, \beta)_{m,n} = 0$$

(c - dp_m p_n)S_R(+, \beta)_{m,n} + (bp_n - ap_m)S_R(-, \beta)_{m,n} = 0,
(8)

where $\beta = +, -, m, n = 1, ..., L$ and a, b, c, d are defined in (A.2). Equations (8) determine the elements of the local matrices $S_R(\alpha, \beta)$ occurring in (5) provided that the determinant of this system of homogeneous linear equations vanishes:

$$(a^{2} + b^{2} - c^{2} - d^{2})p_{m}p_{n} = ab(p_{m}^{2} + pn_{n}^{2}) - cd(1 + p_{m}^{2}p_{n}^{2}).$$
(9)

This determines p_n if p_m is given. Setting

$$p_m = k^{1/2} \operatorname{sn}(u),$$
 (10)

it follows from (A.3) that

$$p_n = k^{1/2} \operatorname{sn}(u \pm 2\eta).$$
 (11)

Baxter selected a solution which has non-vanishing diagonal elements $S_R(\alpha, \beta)_{0,0}$ and $S_R(\alpha, \beta)_{L,L}$. In order to allow $S_R(\alpha, \beta)_{m,n}$ to have non-vanishing diagonal elements $S_R(\alpha, \beta)_{0,0}$ and $S_R(\alpha, \beta)_{L,L}$ equation (9) has to be satisfied for n = m. Then

$$\operatorname{sn}(u) = \operatorname{sn}(u \pm 2\eta). \tag{12}$$

This fixes the parameter u to become $u = K \pm \eta$ and leads to the restriction to discrete η :

$$2L\eta = 2m_1K + \mathrm{i}m_2K'. \tag{13}$$

One obtains from (10) and (11) that

$$p_n = k^{1/2} \operatorname{sn}(K + (2n - 1)\eta)$$
(14)

and from (8)

$$S_{R}(\alpha,\beta)(\nu)_{k,l} = \delta_{k+1,l}u^{\alpha}(\nu+K-2k\eta)\tau_{-k,\beta} + \delta_{k,l+1}u^{\alpha}(\nu+K+2l\eta)\tau_{l,\beta} + \delta_{k,1}\delta_{l,1}u^{\alpha}(\nu+K)\tau_{0,\beta} + \delta_{k,L}\delta_{l,L}u^{\alpha}(\nu+K+2L\eta)\tau_{L,\beta},$$
(15)

for $1 < k \leq L$, $1 < l \leq L$ and where

$$u^{+}(v) = H(v)$$
 $u^{-}(v) = \Theta(v)$ (16)

if

$$\eta = m_1 K / L. \tag{17}$$

 $Q_{R,72}$ is the matrix Q_R defined in (5) with S_R given by (15).

It has been shown in [5] that Q_R based on (15) is singular if m_1 is even and L is odd. In the following subsection we show that an alternative construction leads for these η -values to a regular Q_R -matrix.

1.2. Another Q-matrix

To obtain another solution \hat{S}_R of (8) and (9) we consider the possibility that the elements of $\hat{S}_R(\alpha, \beta)_{m,n}$ form cycles

$$\hat{S}_R(\alpha,\beta)_{1,2}, \hat{S}_R(\alpha,\beta)_{2,3}, \ldots, \hat{S}_R(\alpha,\beta)_{L-1,L}, \hat{S}_R(\alpha,\beta)_{L,1}$$

and

$$\hat{S}_R(\alpha,\beta)_{2,1}, \hat{S}_R(\alpha,\beta)_{3,2}, \dots, \hat{S}_R(\alpha,\beta)_{L,L-1}, \hat{S}_R(\alpha,\beta)_{1,L}$$

instead of imposing the condition that $\hat{S}_R(\alpha, \beta)_{m,n}$ has two diagonal elements. In this case a set of functions p_n consistent with (10) and (11) is

$$p_n = k^{1/2} \operatorname{sn}(t + (2n - 1)\eta).$$
(18)

From the condition that

$$\hat{S}_R(\alpha,\beta)_{L,L+1} = \hat{S}_R(\alpha,\beta)_{L,1} \tag{19}$$

it follows that p_1 and p_L must have arguments which differ by 2η :

$$sn(t + (2L - 1)\eta) = sn(t + \eta - 2\eta).$$
(20)

This is satisfied if

$$2L\eta = 4mK + 2im_2K'.$$
(21)

This condition differs from (13). The solution of equations (8) with the set of p_n -functions (18) as input is

$$\hat{S}_{R}(\alpha,\beta)_{k,l} = \delta_{k+1,l} w^{\alpha} (v-t-2k\eta) \tau_{\beta,-k} + \delta_{k,l+1} u^{\alpha} (v+t+2l\eta) \tau_{\beta,l}$$
(22)

with u^{α} defined in equation (16) and w^{α} is given by

$$w^{+}(v) = -H(v)$$
 $w^{-}(v) = \Theta(v).$ (23)

Note that the first component of w^{α} differs from u^{+} by a minus sign.

We consider only the case $m_2 = 0$ in (21). Then

$$\eta = 2mK/L. \tag{24}$$

We shall denote the Q_R -, Q_L - and Q-matrices derived from \hat{S}_R , \hat{S}_L by \hat{Q}_R , \hat{Q}_L and \hat{Q} . We distinguish the following cases.

- (1) If *L* is odd the resulting \hat{Q} -matrices cover exactly the set of discrete η -values which is missing in the original solution (15)–(16). We note that for t = K this solution becomes identical to case (15)–(16) with singular \hat{Q}_R . But for generic *t* (especially t = 0) \hat{Q}_R is regular. It must be stressed, however, that the regularity has not been proved analytically but numerically for sufficiently large systems to allow the occurrence of degenerate eigenvalues of the transfer matrix *T*. See also appendix C of [1, 5].
- (2) *L* is even but both $L_1 = L/2$ and *m* are odd.

Then $\eta = mK/L_1$ is that set of η -values for which solution (15)–(16) leads to regular Q_R matrices. It turns out that in this case the \hat{Q}_R -matrix resulting from solution (22) is singular.

(3) L and L/2 are even and m is odd.

In this case both solutions (15)–(16) and (22) give regular Q_R -matrices. But the matrices $\hat{S}_R(\alpha, \beta)$ differ in size by a factor of 2.

The conclusion is that the two sets of Q-matrices (15)–(16) and (22) are complementary in the sense that for $\eta = mK/L$ and odd L what is missing in the first set is present in the second and vice versa.

1.3. The matrix \hat{Q}_L

To get finally a Q-matrix which commutes with the transfer matrix T and satisfies equation (1) Baxter introduced a second matrix Q_L . By transposing equation (6) and replacing v by -v one obtains

$$Q_L(v)T(v) = [\rho h(v-\eta)]^N Q_L(v+2\eta) + [\rho h(v+\eta)]^N Q_L(v-2\eta)$$
(25)

with

$$Q_L(v) = Q_R^t(-v) \tag{26}$$

and

$$[Q_L(v)]_{\alpha|\beta} = \operatorname{Tr} S_L(\alpha_1, \beta_1) S_L(\alpha_2, \beta_2) \cdots S_L(\alpha_N, \beta_N).$$
(27)

We perform this construction for the new \hat{Q}_R matrix. The local matrices \hat{S}_L are obtained from (22)

$$\hat{S}_L(\alpha,\beta)_{k,l}(v) = \hat{S}_R(\beta,\alpha)_{k,l}(-v)$$
(28)

$$\hat{S}_{L}(\alpha,\beta)_{k,l} = \delta_{k+1,l}\tau_{\alpha,-k}u^{\beta}(v+t+2k\eta) + \delta_{k,l+1}\tau_{\alpha,l}w^{\beta}(v-t-2l\eta).$$
(29)

1.4. The relation $Q_L(u)Q_R(v) = Q_L(v)Q_R(u)$

To prove that the Q-matrix defined by

$$Q(v) = Q_R(v)Q_R^{-1}(v_0)$$
(30)

commutes with the transfer matrix T Baxter shows in [1] that the relation

$$Q_L(v)Q_R(u) = Q_L(u)Q_R(v)$$
(31)

holds. Then

$$Q(v) = Q_L^{-1}(u)Q_L(v) = Q_R(v)Q_R^{-1}(u)$$
(32)

commutes with T(v). To prove (31) it is shown in [1] that $S_L(\alpha, \gamma)_{m,n}(u)S_R(\gamma, \beta)_{m'n'}(v)$ and $S_L(\alpha, \gamma)_{m,n}(v)S_R(\gamma, \beta)_{m'n'}(u)$ are related by a similarity transformation:

$$S_L(\alpha, \gamma)_{m,n}(u)S_R(\gamma, \beta)_{m'n'}(v) = Y_{m,m';k,k'}S_L(\alpha, \gamma)_{k,l}(v)S_R(\gamma, \beta)_{k'l'}(u)Y_{l,l';n,n'}^{-1},$$
with diagonal matrix Y ,
$$(33)$$

$$Y_{m,m';k,k'} = y_{m,m'} \delta_{m,k} \delta_{m',k'}.$$
(34)

To investigate whether the matrices \hat{Q}_R and \hat{Q}_L defined in (5), (22) and (27), (29) fulfil such a relation we define a series of abbreviations. According to (22) we write

$$\hat{S}_R(\alpha,\beta)_{m,n} = \Phi^{\alpha}_{m,n} \bar{\tau}^{\beta}_{m,n}, \tag{35}$$

where

$$\Phi^{\alpha}_{m,n} = \epsilon^{\alpha}_{m,n} f^{\alpha}(v_{m,n}) \tag{36}$$

$$v_{m,n} = \delta_{m-1,n}(v+t+2n\eta) + \delta_{m+1,n}(v-t-2m\eta)$$
(37)

$$\epsilon_{m,n}^{\alpha} = \delta_{m-1,n} - \alpha \delta_{m+1,n} \qquad \alpha = \pm 1 \tag{38}$$

$$\bar{\tau}_{m,n}^{\beta} = \delta_{m-1,n} \tau_{\beta,n} + \delta_{m+1,n} \tau_{\beta,-m} \tag{39}$$

and $f^+(v) = H(v)$, $f^-(v) = \Theta(v)$, $\delta_{m+L,n} = \delta_{m,n}$. Equivalently, we write following (29):

$$\hat{S}_L(\alpha,\beta)_{m,n} = \tau_{m,n}^{\prime\alpha} \chi_{m,n}^{\beta},\tag{40}$$

where

$$\chi^{\beta}_{m,n} = \lambda^{\beta} f^{\beta}(u_{m,n}) \tag{41}$$

$$u_{m,n} = \delta_{m-1,n}(v - t - 2n\eta) + \delta_{m+1,n}(v + t + 2m\eta)$$
(42)

$$\lambda_{m,n}^{\beta} = -\beta \delta_{m-1,n} + \delta_{m+1,n} \tag{43}$$

$$\bar{\tau'}_{m,n}^{\alpha} = \delta_{m-1,n} \tau'_{\alpha,n} + \delta_{m+1,n} \tau'_{\alpha,-m}.$$
(44)

It follows then from (35) and (40)

$$\hat{S}_{L}(\alpha,\gamma)_{m,n}(u)\hat{S}_{R}(\gamma,\beta)_{m'n'}(v) = \tau_{m,n}^{\prime\alpha}\chi_{m,n}^{\gamma}(u)\Phi_{m',n'}^{\gamma}(v)\tau_{m',n'}^{\beta}$$
(45)

and from (36) and (41) one obtains

$$\chi_{m,n}^{\gamma}(u)\Phi_{m',n'}^{\gamma}(v) = (\delta_{m+1,n}\delta_{m'+1,n'} + \delta_{m-1,n}\delta_{m'-1,n'})(\Theta(u_{m,n})\Theta(v_{m',n'}) - H(u_{m,n})H(v_{m',n'})) + (\delta_{m+1,n}\delta_{m'-1,n'} + \delta_{m-1,n}\delta_{m'+1,n'})(\Theta(u_{m,n})\Theta(v_{m',n'}) + H(u_{m,n})H(v_{m',n'})),$$
(46)

with non-vanishing elements

$$\chi_{m,m+1}^{\gamma}(u)\Phi_{m',m'+1}^{\gamma}(v) = \Theta(u_{m,m+1})\Theta(v_{m',m'+1}) - H(u_{m,m+1})H(v_{m',m'+1})$$
(47)

$$\chi_{m,m-1}^{\gamma}(u)\Phi_{m',m'-1}^{\gamma}(v) = \Theta(u_{m,m-1})\Theta(v_{m',m'-1}) - H(u_{m,m-1})H(v_{m',m'-1})$$
(48)

$$\chi_{m,m+1}^{\gamma}(u)\Phi_{m',m'-1}^{\gamma}(v) = \Theta(u_{m,m+1})\Theta(v_{m',m'-1}) + H(u_{m,m+1})H(v_{m',m'-1})$$
(49)

$$\chi_{m,m-1}^{\gamma}(u)\Phi_{m',m'+1}^{\gamma}(v) = \Theta(u_{m,m-1})\Theta(v_{m',m'+1}) + H(u_{m,m-1})H(v_{m',m'+1}).$$
(50)

The arguments are

$$u_{m,m+1}(u) - v_{m',m'+1}(v) = u - v + 2(m + m')\eta + 2t$$

$$u_{m,m-1}(u) - v_{m',m'-1}(v) = u - v - 2(n + n')\eta - 2t$$

$$u_{m,m+1}(u) - v_{m',m'-1}(v) = u - v + 2(m - m' + 1)\eta$$

$$u_{m,m-1}(u) - v_{m',m'+1}(v) = u - v - 2(m - m' - 1)\eta$$

$$u_{m,m+1}(u) + v_{m',m'+1}(v) = u + v + 2(m - m')\eta$$

$$u_{m,m-1}(u) + v_{m',m'-1}(v) = u + v + 2(-n + n')\eta$$

$$u_{m,m+1}(u) + v_{m',m'-1}(v) = u + v + 2(m + n')\eta + 2t$$

$$u_{m,m-1}(u) + v_{m',m'+1}(v) = u + v - 2(n + m')\eta - 2t.$$

(51)

To rewrite (47)–(50), we use

$$\Theta(u)\Theta(v) + H(u)H(v) = cf_{+}(u+v)g_{+}(u-v)$$
(52)

$$\Theta(u)\Theta(v) - H(u)H(v) = cf_{-}(u+v)g_{-}(u-v)$$
(53)

$$f_{+}(u) = H((iK'+u)/2)H((iK'-u)/2)g_{+}(u) = H_{1}((iK'+u)/2)H_{1}((iK'-u)/2)$$
(54)

$$f_{-}(u) = H_{1}((iK'+u)/2)H_{1}((iK'-u)/2)g_{-}(u) = H((iK'+u)/2)H((iK'-u)/2).$$
 (55)

We need especially the following properties of g_{\pm} :

$$g_{\pm}(-u) = g_{\pm}(u)$$
 $g_{\pm}(u+4K) = g_{\pm}(u).$ (56)

After insertion of (52)–(55) into (47)–(50) we get

$$\chi_{m,m+1}^{\gamma}(u)\Phi_{m',m'+1}^{\gamma}(v) = cf_{-}(u+v+2(m-m')\eta)g_{-}(u-v+2(m+m')\eta+2t)$$
(57)

$$\chi_{m,m-1}^{\gamma}(u)\Phi_{m',m'-1}^{\gamma}(v) = cf_{-}(u+v+2(m'-m)\eta)g_{-}(u-v-2(n+n')\eta-2t)$$
(58)

$$\chi_{m,m+1}^{\gamma}(u)\Phi_{m',m'-1}^{\gamma}(v) = cf_{+}(u+v+2(m+n')\eta+2t)g_{+}(u-v+2(m-m'+1)\eta)$$
(59)

$$\chi_{m,m-1}^{\gamma}(u)\Phi_{m',m'+1}^{\gamma}(v) = cf_{+}(u+v-2(n+m')\eta-2t)\eta)g_{+}(u-v-2(m-m'-1)\eta).$$
(60)

It now remains to show that a $L^2 \times L^2$ matrix Y exists such that equation (33) is satisfied for \hat{S}_R and \hat{S}_L . As τ and τ' occurring in the definition of \hat{S}_R and \hat{S}_L are free parameters we obtain from (33)

$$\chi_{m,n}^{\gamma}(u)\Phi_{m',n'}^{\gamma}(v) = Y_{m,m';k,k'}\chi^{\gamma}(v)_{k,l}\Phi^{\gamma}(u)_{k',l'}Y_{l,l';n,n'}^{-1}.$$
(61)

Taking tentatively Y to be diagonal

$$Y_{m,m';k,k'} = y_{m,m'} \delta_{m,k} \delta_{m',k'},$$
(62)

we get

$$\chi_{m,n}^{\gamma}(u)\Phi_{m',n'}^{\gamma}(v) = \frac{y_{m,m'}}{y_{n,n'}}\chi_{m,n}^{\gamma}(v)\Phi_{m',n'}^{\gamma}(u)$$
(63)

and it follows from (57)–(60)

$$y_{m+1,m'+1} = y_{m,m'} \frac{g_{-}(u - v - 2(m + m')\eta - 2t)}{g_{-}(u - v + 2(m + m')\eta + 2t)}$$
(64)

$$y_{m+1,m'-1} = y_{m,m'} \frac{g_+(u-v-2(m-m'+1)\eta)}{g_+(u-v+2(m-m'+1)\eta))}.$$
(65)

To prove that a matrix Y can be found such that (61) is satisfied we have to show that the set of equations (64)–(65) is free from contradictions on the torus of size $L \times L$ where

$$y_{m+L,n+L} = y_{m,n}.$$
 (66)

It follows from equation (64) that

$$y_{m+L,n+L} = \frac{g_{-}(u-v-2(m+n)\eta - 4(L-1)\eta - 2t)}{g_{-}(u-v+2(m+n)\eta + 4(L-1)\eta + 2t)} \times \frac{g_{-}(u-v-2(m+n)\eta - 4(L-2)\eta - 2t)}{g_{-}(u-v+2(m+n)\eta + 4(L-2)\eta + 2t)} \cdots \frac{g_{-}(u-v-2(m+n)\eta - 2t)}{g_{-}(u-v+2(m+n)\eta + 2t)} y_{m,n}.$$
(67)

The factor $g_{-}(u - v - 2(m + n)\eta + 4r_2\eta - 2t)$ in the numerator cancels the factor $g_{-}(u - v + 2(m + n)\eta + 4r_1\eta + 2t)$ in the denominator if t = 0 and

$$-2(m+n)\eta - 4r_2\eta = 2(m+n)\eta + 4r_1\eta + 4kK$$
(68)

for arbitrary k and if we set $k = 2m_1k_1$ for integer k_1 :

$$r_2 = k_1 L - m - n - r_1. ag{69}$$

It follows that for each factor in the numerator of equation (67) there is a factor in the denominator against which it cancels. Similarly we derive from equation (65) that

$$y_{m,n} = y_{m-L,n+L}.$$
 (70)

We have shown that all $y_{m,n}$ can be determined from a single element (e.g. $y_{1,1}$) consistently if t = 0. This conclusion cannot be drawn for $t \neq 0$. A numerical test of (31) shows that it is not satisfied for $t \neq 0$, and therefore no similarity transformation (33) exists for $t \neq 0$.

We summarize what has been found in this section.

We have attained our goal to construct a *Q*-matrix which exists for $\eta = 2mK/L$ for odd *L*:

The \hat{Q}_R -matrix defined in equation (5) with local matrices \hat{S}_R defined in (22) is regular.

If the parameter t is set to zero relation (31) is satisfied.

Then $\hat{Q}(v) = \hat{Q}_R(v)\hat{Q}_R^{-1}(v_0)$ satisfies equation (1) and commutes with the transfer matrix *T*.

2. Quasiperiodicity properties of Q

It is easily seen that $Q_{72,R}(v)$ and $\hat{Q}_{R}(v, t)$ satisfy

$$\hat{Q}_{72,R}(v+2K) = S\hat{Q}_{72,R}(v) \qquad \hat{Q}_{R}(v+2K,t) = S\hat{Q}_{R}(v,t).$$
(71)

It is of great importance to find the quasiperiodicity properties of the *Q*-matrices in the complex *v*-plane. We do that in this section for $Q_{72,R}(v)$ and $\hat{Q}_R(v, t = 0)$. It is well known that the quasiperiod of Q_{73} is iK'. See [9] for details. We shall present plausibility arguments for the statement that $Q_{72,R}$ as well as \hat{Q} are singular matrices if their quasiperiod is iK'.

2.1. Quasiperiodicity properties of Q_{72}

We get from equations (15), (A.4), (A.5) and from

$$\eta = m_1 K / L \tag{72}$$

the relations

$$S_{R}(\pm,\beta)_{k,k+1}(v+i\mathsf{K}') = f(v)\exp(+i\pi k\eta/K)S_{R}(\mp,\beta)(v)_{k,k+1}$$

$$S_{R}(\pm,\beta)_{k+1,k}(v+i\mathsf{K}') = f(v)\exp(-i\pi k\eta/K)S_{R}(\mp,\beta)(v)_{k+1,k}$$

$$S_{R}(\pm,\beta)_{1,1}(v+i\mathsf{K}') = f(v)S_{R}(\mp,\beta)(v)_{1,1}$$

$$S_{R}(\pm,\beta)_{1,1}(v+i\mathsf{K}') = f(v)S_{R}(\mp,\beta)(v)_{1,1}$$
(73)

 $\mathsf{S}_R(\pm,\beta)_{L,L}(v+\mathsf{i}\mathsf{K}')=(-1)^{m_1}f(v)\mathsf{S}_R(\mp,\beta)(v)_{L,L},$

where

$$f(v) = q^{-1/4} \exp\left(-\frac{\mathrm{i}\pi v}{2K}\right). \tag{74}$$

The similarity transformation

$$\bar{S}(\alpha,\beta)_{i,l} = A_{i,j}\hat{S}(\alpha,\beta)_{j,k}A_{k,l}^{-1},\tag{75}$$

with

$$A_{k,l} = \delta_{k,l} \exp\left(\frac{\mathrm{i}\pi}{2K}(k-1)k\eta\right) a_1 \tag{76}$$

leads to

$$\begin{aligned} \tilde{S}_{R}(\pm,\beta)_{k,k+1}(v+i\mathsf{K}') &= f(v)S_{R}(\mp,\beta)(v)_{k,k+1} \\ \tilde{S}_{R}(\pm,\beta)_{k+1,k}(v+i\mathsf{K}') &= f(v)S_{R}(\mp,\beta)(v)_{k+1,k} \\ \tilde{S}_{R}(\pm,\beta)_{1,1}(v+i\mathsf{K}') &= f(v)S_{R}(\mp,\beta)(v)_{1,1} \\ \tilde{S}_{R}(\pm,\beta)_{L,L}(v+i\mathsf{K}') &= (-1)^{m_{1}}f(v)S_{R}(\mp,\beta)(v)_{L,L}.
\end{aligned} \tag{77}$$

If m_1 is even it follows that

$$\tilde{\mathsf{S}}_{R}(\alpha,\beta)_{k,l}(\nu+\mathsf{i}\mathsf{K}') = f(\nu)R(\alpha,\gamma)\mathsf{S}_{R}(\gamma,\beta)(\nu)_{k,l},\tag{78}$$

where R is defined in equation (4) and

$$Q_{R,72}(v + i\mathsf{K}') = f(v)^N R Q_{R,72}(v).$$
(79)

However, it is well known [5] that for even $m_1 Q_{R,72}(v)$ is singular and consequently relation (79) cannot be upgraded from $Q_{R,72}$ to Q_{72} . It is shown in [5] that instead of (79) the following relation holds:

$$Q_{R,72}(v+2iK') = q^{-N} \exp(-iN\pi v/K) Q_{R,72}(v),$$
(80)

which is correct for all $\eta = m_1 K/L$. Provided m_1 is odd it follows

$$Q_{72}(v+2iK') = q^{-N} \exp(-iN\pi v/K) Q_{72}(v).$$
(81)

2.2. Quasiperiodicity properties of \hat{Q}

We obtain from equation (22)

$$\hat{S}_{R}(\pm,\beta)(v+i\mathbf{K}')_{k,k+1} = f(v)(-i)\exp(+i\pi k\eta/K)\hat{S}_{R}(\mp,\beta)(v)_{k,k+1}$$
(82)

$$\hat{S}_{R}(\pm,\beta)(v+i\mathsf{K}')_{k+1,k} = f(v)(\pm i)\exp(-i\pi k\eta/K)\hat{S}_{R}(\mp,\beta)(v)_{k+1,k}.$$
(83)

Perform the similarity transformation

$$\bar{S}(\alpha,\beta)_{i,l} = A_{i,j}\hat{S}(\alpha,\beta)_{j,k}A_{k,l}^{-1},\tag{84}$$

with

$$A_{k,l} = \delta_{k,l} (-\mathbf{i})^{k-1} \exp\left(\frac{\mathbf{i}\pi}{2K} (k-1)k\eta\right) a_1.$$
(85)

Then for k < L,

$$\bar{S}_{R}(\pm,\beta)(v+i\mathsf{K}')_{k,k+1} = f(v)\hat{S}_{R}(\mp,\beta)(v)_{k,k+1}$$
(86)

$$\bar{S}_{R}(\pm,\beta)(v+i\mathsf{K}')_{k+1,k} = f(v)\hat{S}_{R}(\mp,\beta)(v)_{k+1,k}$$
(87)

and for k = L,

$$\bar{S}_{R}(\pm,\beta)(v+\mathsf{i}\mathsf{K}')_{L,1} = f(v)\exp\left[\frac{+\mathsf{i}\pi}{2}((2m_{1}-1)L+2m_{1})\right]\hat{S}_{R}(\mp,\beta)(v)_{k,k+1}$$
(88)

$$\bar{S}_{R}(\pm,\beta)(v+\mathsf{i}\mathsf{K}')_{1,L} = f(v)\exp\left[\frac{-\mathsf{i}\pi}{2}((2m_{1}-1)L+2m_{1})\right]\hat{S}_{R}(\mp,\beta)(v)_{k+1,k}.$$
(89)

We find that if

$$\exp\left[\frac{-i\pi}{2}((2m_1 - 1)L + 2m_1)\right] = 1,$$
(90)

 \hat{Q}_R satisfies the relation

$$\hat{Q}_R(\nu + i\mathsf{K}') = q^{-N/4} \exp\left(-\frac{i\pi N\nu}{2K}\right) R\hat{Q}_R(\nu)$$
(91)

which is the same as (79) for $Q_{R,72}$. This happens only for

- (I) even m_1 if $L = 4 \times$ integer
- (II) odd m_1 if $L = 2 \times$ odd integer.

These are exactly those cases in which \hat{Q}_R is singular. Like equation (79) equation (91) does not give the corresponding relation for the Q-matrix \hat{Q} .

In the following paragraph Q denotes either Q_{72} or \hat{Q} .

We note that if the relation

$$Q(v + iK') = q^{-N/4} \exp\left(-\frac{i\pi Nv}{2K}\right) RQ(v)$$
(92)

were correct then it would follow that

$$q(v + i\mathbf{K}')|q\rangle = q^{-N/4} \exp\left(-\frac{i\pi Nv}{2K}\right) q(v)R|q\rangle, \tag{93}$$

where $|q\rangle$ denotes an arbitrary eigenvector of Q(v) and q(v) is its eigenvalue.

In other words: all eigenvectors of $\hat{Q}(v)$ would be eigenvectors of R. It is however well known [5] that the eigenvectors of $Q_{72}(v)$ which are degenerate eigenvectors of the transfer matrix T are generally not eigenvectors of R.

Equations (79) and (91) allow a coherent explanation of the fact that Q_R is singular for one set of η values and regular for another. Under the assumption that if Q exists there are eigenstates of Q which are not eigenstates of R, Q_R cannot be regular if in case of Q_{72} , m_1 is even or in case of \hat{Q} (90) is satisfied. This also explains naturally the observation that for fixed L and sufficiently small NQ_R^{-1} exists always as then all states are singlets and (79) and (91) do not lead to contradictions when upgraded from Q_R to Q.

Using the method used in this section it can easily be shown that always

$$\hat{Q}_{R}(v+2iK') = q^{-N} \exp(-iN\pi v/K) \hat{Q}_{R}(v)$$
(94)

and consequently if \hat{Q}_R^{-1} exists

$$\hat{Q}(v+2iK') = q^{-N} \exp(-iN\pi v/K)\hat{Q}(v).$$
(95)

3. The properties of \hat{Q} for t = 0

It follows from (95) that as shown for Q_{72} in [5], $\hat{Q}(v)$ may be written as

$$\hat{Q}(v) = \hat{\mathcal{K}}(q; v_k) \exp(i(n_B - v)\pi v/2K) \prod_{j=1}^{n_B} h(v - v_j^B) \times \prod_{j=1}^{n_L} H(v - iw_j) H(v - iw_j - 2\eta) \cdots H(v - iw_j - 2(L - 1)\eta)$$
(96)
$$2n_B + Ln_L = N.$$
(97)

$$2n_B + Ln_L = N.$$

 n_B is the number of Bethe roots v_k^B and n_L is the number of exact Q-strings of length L. The sum rules (2.16)–(2.21) of [5] are also true for $\hat{Q}(v)$.

From (71) and

$$\hat{Q}_L(v)\hat{Q}_R(u) = \hat{Q}_L(u)\hat{Q}_R(v)$$
 (98)

follows that

$$[S, \hat{Q}(v)] = 0. (99)$$

Finally, we find numerically that the functional relation (3) which was originally conjectured in [5] is also satisfied for $\hat{Q}(v)$.

4. The matrix $\hat{Q}(v,t)$

We have shown that $\hat{Q}_R(v)$ satisfies the $T - \hat{Q}_R$ relation and found that it is not singular. But the proof of relation (31) failed for parameter $t \neq 0$. One finds numerically that (31) is in fact violated for systems large enough to allow degenerate eigenvalues of the transfer matrix. Therefore, the question arises whether $\hat{Q}_R(v, t)$ is useful at all. Surprisingly, we find numerically that for $\eta = 2mK/L$ and odd L despite

$$\hat{Q}_{L}^{-1}(v_{0},t)\hat{Q}_{L}(v,t) \neq \hat{Q}_{R}(v,t)\hat{Q}_{R}^{-1}(v_{0},t)$$
(100)

both matrices

$$\hat{Q}^{(L)}(v,t) = \hat{Q}_{L}^{-1}(v_{0},t)\hat{Q}_{L}(v,t) \quad \text{and} \quad \hat{Q}^{(R)}(v,t) = \hat{Q}_{R}(v,t)\hat{Q}_{R}^{-1}(v_{0},t) \quad (101)$$

commute with the transfer matrix *T*. Furthermore, we find that in the cases studied $\hat{Q}^{(L)}(v, t)$ and $\hat{Q}^{(R)}(v, t)$ have the same eigenvalues. This means that there exists a matrix *A* with $\hat{Q}^{(L)}(v, t) = A\hat{Q}^{(R)}(v, t)A^{-1}$ and consequently instead of (31)

$$\hat{Q}_{L}(v)A\hat{Q}_{R}(u) = \hat{Q}_{L}(u)A\hat{Q}_{R}(v)$$
(102)

should hold. A consequence of (100) is that

$$[\hat{Q}^{(R)}(v_1, t), \hat{Q}^{(R)}(v_2, t)] \neq 0$$
(103)

as one needs (31) to prove that *Q*-matrices with different arguments commute (see 9.48.41 in [9]). We find that like Q_{72} the matrix \hat{Q} does not commute with *R*:

$$[R, \hat{Q}(v, t)] \neq 0, \tag{104}$$

but unlike Q_{72} as a consequence of (100) does also not commute with S for $t \neq 0$:

$$[S, \hat{Q}(v, t)] \neq 0.$$
 (105)

This is possible because the degenerate subspaces of T have elements with both eigenvalues $\nu' = 0, 1$ of S if $\eta = 2m_1 K/L$ and L is odd.

These properties of $\hat{Q}^{(L)}(v, t)$ and $\hat{Q}^{(R)}(v, t)$ imply that they act as non-Abelian symmetry operators in all degenerate subspaces of the set of eigenvectors of *T*. We finally mention that whereas the six-vertex limit of $Q_{R,72}$ does not exist it exists for \hat{Q}_R . The limit of $\hat{Q}_R(v, t = iK'/2)$ for elliptic nome $q \to 0$ is well defined. Using

$$\lim_{q \to 0} H(u \pm iK'/2) = \exp(\mp i(u - \pi/2)) \lim_{q \to 0} \Theta(u \pm iK'/2) = 1,$$
(106)

one gets a regular limiting \hat{Q}_R -matrix. It has been checked numerically that the resulting \hat{Q} -matrix commutes with T_{6v} .

Acknowledgment

I am pleased to thank Professor Barry M McCoy for helpful comments and suggestions.

Appendix

The transfer matrix of the eight-vertex model is

$$T(v)|_{\mu,\nu} = \operatorname{Tr} W_8(\mu_1, \nu_1) W_8(\mu_2, \nu_2) \cdots W_8(\mu_N, \nu_N),$$
(A.1)

where in the conventions of (6.2) of [1]

$$\begin{split} & W_8(1,1)|_{1,1} = W_8(-1,-1)|_{-1,-1} = a = \rho \Theta(2\eta)\Theta(v-\eta)H(v+\eta) \\ & W_8(-1,-1)|_{1,1} = W_8(1,1)|_{-1,-1} = b = \rho \Theta(2\eta)H(v-\eta)\Theta(v+\eta) \\ & W_8(-1,1)|_{1,-1} = W_8(1,-1)|_{-1,1} = c = \rho H(2\eta)\Theta(v-\eta)\Theta(v+\eta) \\ & W_8(1,-1)|_{1,-1} = W_8(-1,1)|_{-1,1} = d = \rho H(2\eta)H(v-\eta)H(v+\eta). \end{split}$$
(A.2)

Relations used in the text. See e.g. [10]

$$sn(u - v) = \frac{sn(u) cn(v) dn(v) - sn(v) cn(u) dn(u)}{1 - k^2 sn^2(u) sn^2(v)}$$
(A.3)

$$H(v + iK') = iq^{-1/4} \exp\left(-\frac{i\pi v}{2K}\right)\Theta(v)$$
(A.4)

$$\Theta(v + iK') = iq^{-1/4} \exp\left(-\frac{i\pi v}{2K}\right) H(v).$$
(A.5)

References

- [1] Baxter R J 1972 Partition function of the eight vertex model Ann. Phys. 70 193
- Baxter R J 1973 Eight-vertex model in lattice statistics and one-dimensional anisotropic Heisenberg chain: I. Some fundamental eigenvectors Ann. Phys. 76 1
- [3] Baxter R J 1973 Eight-vertex model in lattice statistics and one-dimensional anisotropic Heisenberg chain: II. Equivalence to a generalized ice-type lattice model Ann. Phys. 76 25
- [4] Baxter R J 1973 Eight-vertex model in lattice statistics and one-dimensional anisotropic Heisenberg chain: III. Eigenvectors of the transfer matrix and the Hamiltonian Ann. Phys. 76 48
- [5] Fabricius K and McCoy B M 2003 New developments in the eight vertex model J. Stat. Phys. 111 323–37
- [6] Fabricius K and McCoy B M 2004 Functional equations and fusion matrices for the eight-vertex model *Publ*. *RIMS* 40 905
- [7] Fabricius K and McCoy B 2002 Evaluation parameters and Bethe roots for the six-vertex model at roots of unity *MathPhys Odyssey 2001 (Progress in Mathematical Physics 23)* ed M Kashiwara and T Miwa (Boston: Birkhäuser) pp 119–44
- [8] Bazhanov V V and Stroganov Yu G 1990 Chiral Potts Model as a Descendant of the Six-vertex Model J. Stat. Phys. 59 799
- [9] Baxter R J 1982 Exactly Solved Models (London: Academic)
- [10] Whittaker E T and Watson G N A Course of Modern Analysis (Cambridge: Cambridge University Press)
- [11] Roan Sih-shyr 2006 Preprint cond-mat/0611316